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Remark on the semiconductor equations (**)

1 - Introduction

We consider the Van Roosbroeck model for the carrier transport in a semiconductor device [19]; such model describes the transport of carriers in absence of ionized impurities by means of the following system of non-linear partial differential equations:

$$(1.1) \hspace{1cm} \begin{aligned} p_t - \nabla \cdot (D_1 \nabla p + \mu_1 p \nabla \Phi) &= R(p, \, n) \\ n_t - \nabla \cdot (D_2 \nabla n - \mu_2 n \nabla \Phi) &= R(p, \, n) \quad \text{on } Q_T \\ - \nabla \cdot (\varepsilon \nabla \Phi) &= p - n \end{aligned}$$

where p and n are the densities of mobile holes and electrons, while Φ is the electric potential. The positive constants D_i , μ_i and ε are respectively the diffusion coefficients, the mobilities of holes and electrons, and the dielectric permittivity. The term R(p,n) represents the law of recombination which, according to the Shockley-Read-Hall model, can be expressed by

$$R(p, n) = \frac{1 - pn}{r_0 + r_1 p + r_2 n}$$

where r_i , i = 0, 1, 2, are positive constants. The domain Q_T is the product $\Omega \times (0, T)$, where Ω is an open bounded subset of \mathbb{R}^N with regular boundary.

Due to their physical and technological interest (see e.g. [11] and [15] for more detailed information on the physical background), the above equations have recently received a great deal of attention from the mathematical point of view. The

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problems of existence and uniqueness, regularity, and asymptotic behaviour of the solutions have been widely studied by different authors, as M. S. Mock [12], [13], [14], T. I. Seidman [16], [17], H. Beirão da Veiga [2], [3], [4], H. Gajewski [7], [8], and K. Gröger [7], [8], [9], [10].

In this paper we show the existence of an invariant region for the mobile carriers p and n. That immediately yields the existence of a solution global in time to (1.1), as well as an a priori sup-norm bound for p and n. The main aim of this paper is then to notice that such results, which have been already proved in the literature by using essentially a more complicated version of the maximum principle [17], [9], can be obtained in a quite simplified way.

In Section 2 we briefly recall the definition of invariant region, and a basic related criterion, while in Section 3 we state and prove the existence of an invariant region for (1.1).

2 - Invariant regions for parabolic systems

Let us introduce the notion of invariant region and some related results [6]. For a complete survey about invariant regions, we refer to the book by J. Smoller [18].

Let us consider the weakly coupled parabolic system:

(2.1)
$$(u_k)_t = D \Delta u_k + \sum_{j=1}^N a_j(x, t)(u_k)_{x_j} + \psi_k(u_1, \dots, u_M) \quad \text{on } Q_T$$

$$u_k(x, 0) = (u_k)_0(x) \quad \text{on } \Omega$$

$$u_k(x, t) = (u_k)_k(x, t) \quad \text{on } \partial \Omega \times (0, T)$$

where k = 1, ..., M. Here D denotes a positive constant, while the coefficients a_j , the vector field $V = (\psi_1, ..., \psi_M)$, as well as the initial and boundary data $(u_k)_0$ and $(u_k)_b$ are assumed to be continuous. In the following \boldsymbol{u} will be a vector of components $(u_1, ..., u_M)$.

Definition 1. A subset S of \mathbb{R}^M is said to be an *invariant region* for the parabolic system (2.1) when the following condition holds: if the initial and boundary data of (2.1) lie in S, i.e. if

$$u_0(x) \in S \quad \forall x \in \Omega \quad \text{and} \quad u_b(x, t) \in S \quad \forall (x, t) \in \partial \Omega \times (0, T)$$

then every solution u(x, t) of (2.1) belonging to $C^2(Q_T) \cap C^0(\overline{Q}_T)$ remains in S, i.e. $u(x, t) \in S$, $\forall (x, t) \in Q_T$.

We recall the following result [20], [1], which is the main tool used in the proof of Theorem 2.

Theorem 1. Let S be a closed and convex subset of \mathbb{R}^M such that the vector field V does not point out of S whenever \mathbf{u} is on the boundary of S, that is

$$(2.2) V(u) \cdot v \leq 0$$

for every point $u \in \partial S$ and for every v outward normal vector at u. Then S is an invariant region for the parabolic system (2.1).

Remark 1.

- i. Theorem 1 holds even if ∂S contains singular points. In this case we take as outward normal at a singular point u_s any vector \mathbf{v} of \mathbf{R}^M such that S is contained in the halfspace with exterior normal \mathbf{v} ; the existence of such a vector \mathbf{v} is guaranteed by the hypothesis on S of being a closed and convex set [5].
- ii. Let α_k and β_k be real numbers, for $k=1,\ldots,M$, and suppose that the set $S \subset \mathbb{R}^M$ is of the type

$$S = \{ \boldsymbol{u} \in \boldsymbol{R}^{M} \colon \alpha_{k} \leq u_{k} \leq \beta_{k}, \quad k = 1, \dots, M \}.$$

Then Theorem 1 still holds if we replace the coefficient D in equation (2.1)₁ by a positive constant D_k , depending on the index k.

3 - Existence of an invariant region for the semiconductor equations

We want to use Theorem 1 to prove the existence of an invariant region for the mobile carriers in semiconductor equations.

To this aim we observe that the model (1.1) can be reduced to a weakly coupled parabolic system of type (2.1). Indeed, if we substitute equation $(1.1)_3$ into equations $(1.1)_1$ and $(1.1)_2$, we get:

$$\begin{split} p_t &= D_1 \Delta p + \mu_1 \nabla p \cdot \nabla \Phi - \frac{\mu_1}{\varepsilon} p(p-n) + R(p, n) \\ n_t &= D_2 \Delta n - \mu_2 \nabla n \cdot \nabla \Phi + \frac{\mu_2}{\varepsilon} n(p-n) + R(p, n) \,. \end{split}$$

We consider the above equations coupled with initial data

$$p(x, 0) = p_0(x)$$
 $n(x, 0) = n_0(x)$ on Ω

and with Dirichlet boundary data

$$p = p_b(x, t)$$
 $n = n_b(x, t)$ on $\partial \Omega \times [0, T)$.

For simplicity we suppose that p_0 , n_0 , p_b , n_b are regular functions, satisfying the compatibility condition

$$p_h(x, 0) = p_0(x)$$
 $n_h(x, 0) = n_0(x)$ on Ω .

We are now in a position to state our main result.

Theorem 2. For c sufficiently large, the set

$$Q(c) = \{(p, n) \in \mathbb{R}^2 : 0 \le p \le 1 + c, 0 \le n \le 1 + c\}$$

is an invariant region for the mobile carriers in the semiconductor equations.

Proof. It is sufficient to prove that, for c large enough, the vector field V = (f, g) of components

$$f = -\frac{\mu_1}{\varepsilon} p(p-n) + R(p, n)$$
 $g = \frac{\mu_2}{\varepsilon} n(p-n) + R(p, n)$

satisfies hypothesis (2.2) on the boundary of Q(c). We set

$$\partial \big(Q(c)\big) = \varGamma_1 \cup \varGamma_2 \cup \varGamma_3 \cup \varGamma_4 \ ,$$

where:

$$\begin{split} & \varGamma_1 = \big\{ p = 0 \,,\, 0 \leqslant n \leqslant 1 + c \big\} & \qquad \varGamma_2 = \big\{ n = 0 \,,\, 0 \leqslant p \leqslant 1 + c \big\} \\ & \qquad \varGamma_3 = \big\{ p = 1 + c \,,\, 0 \leqslant n \leqslant 1 + c \big\} & \qquad \varGamma_4 = \big\{ n = 1 + c \,,\, 0 \leqslant p \leqslant 1 + c \big\} \,. \end{split}$$

On Γ_1 we have

$$(V \cdot v)_{|r_1} = -f(0, n) = -\frac{1}{r_0 + r_2 n} < 0$$
 $\forall n \in [0, 1 + c].$

Similarly we obtain $(V \cdot v)_{|_{\Gamma_2}} < 0$. Let us consider Γ_3 .

$$(V \cdot v)_{\mid \Gamma_3} = f(1+c, n) = -\frac{\mu_1}{\varepsilon} (1+c)[(1+c) - n] + \frac{1 - (1+c) n}{r_0 + r_1(1+c) + r_2 n}$$

for any $n \in [0, 1+c]$.

Since

$$f(1+c, 0) = -\frac{\mu_1}{\varepsilon} (1+c)^2 + \frac{1}{r_0 + r_1(1+c)}$$

and

$$f(1+c, 1+c) = \frac{1-(1+c)^2}{r_0 + (r_1 + r_2)(1+c)}$$

are both negative for c large enough, we have only to show that there exists $\bar{c} > 0$ such that the derivative of f(1+c, n) with respect to n is non-negative when c is larger than \bar{c} . Actually, for $n \in [0, 1+c]$, we have the inequality

$$\frac{\partial f}{\partial n}(1+c, n) \ge \frac{\mu_1}{\varepsilon}(1+c) - \frac{(1+c)}{r_0 + r_1(1+c)} - \frac{r_2}{[r_0 + r_1(1+c)]^2}$$

so that

$$\lim_{c \to +\infty} \frac{\partial f}{\partial n} (1 + c, n) = +\infty,$$

which gives the required condition on Γ_3 . In an analogous way one can show that, for c large enough, $(V \cdot \nu)_{|_{\Gamma_4}} < 0$, and the proof is achieved.

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Sommario

Si prova l'esistenza di una regione invariante per il sistema parabolico che descrive il trasporto di lacune ed elettroni nei semiconduttori. Conseguenza immediata di questo risultato è che l'esistenza di una soluzione globale per le equazioni dei semiconduttori e la sua limitatezza nella norma L^{∞} possono essere dimostrate in modo molto semplificato.