HAROLD EXTON (*)

On the orthogonality of the Bessel polynomials (**)

Recently, Srivastava [5] has pointed out that the formula

(1)
$$\int_{0}^{\infty} x^{1-a} e^{-x} y_{m}(x; a, 1) y_{n}(x; a, 1) dx$$

given by Hamza [3] is incorrect. The Bessel polynomial $y_n(x; a, b)$ is given by

(2)
$$y_n(x; a, b) = 2F_0(-n, n+a-1; -; -x/b).$$

In this note, the scaling factor b has been put equal to unity without loss of generality. The Bessel polynomials were first studied on a systematic basis by Krall and Frink [4] who gave the orthogonality relation

(3)
$$\int_{|x|=1} y_m(x;a,1) y_n(x;a,1) w(x) dx = \frac{(-1)^{n+1} n! \Gamma(a)}{(a+2n-1)\Gamma(a+n-1)} \delta_{m,n}$$

where

$$w(x) = \sum_{m=0}^{\infty} \frac{\Gamma(a)}{\Gamma(a+m-1)} \left(-\frac{1}{z}\right)^{m}.$$

Now, the Bessel polynomials $y_n(x; a, 1)$ satisfy the differential equation

(4)
$$x^2y'' + (ax+1)y' - n(n+a-1) y = 0,$$

which may be written in the self-adjoint form

(5)
$$[x^a \exp[-1/x]y']' - n(n+a-1)x^{a-2} \exp[-1/x]y = 0.$$

^(*) Indirizzo: «Nyuggel», Lunabister, Dunrossness, Shetland ZE2 9JH, United Kingdom.

^(**) Ricevuto: 12-VI-1985.

In view of the interest shown in the Bessel polynomials, see Grosswald [2] for example, it seems worthwhile to investigate the matter of their orthogonality further.

From the general theory of Sturm-Liouville systems, the form of (5) leads us to consider integrals of the type

(6)
$$\int_{C} x^{a-2} \exp\left[-\frac{1}{x}\right] y_{m}(x; a, 1) y_{n}(x; a, 1) dx$$

from the point of view of the orthogonality of these polynomials, always provided that the integral (6) converges. The contour C is to be so selected that

(7)
$$[x^a \exp[-1/x]]_C = 0.$$

Hence, this contour may of a simple loop beginning at the origin, encircling the point at $+\infty$ in the positive direction and returning to the origin. If Re(a) < 0, C may also be taken to be a simple path joining 0 and $+\infty$.

Consider the integral

(8)
$$I_{m,n} = \int_{0}^{(\infty+)} x^{a-2} \exp\left[-\frac{1}{x}\right] y_m(x;a,1) y_n(x;a,1) dx.$$

Replace each Bessel polynomial by its hypergeometric series representation (2) and integrate term by term. Hence

(9)
$$I_{m,n} = \sum_{r,s=0}^{m,n} \frac{(-m)_r (-n)_s (m+a-1)_r (n+a-1)_s}{r! \, s!} (-1)^{r+s} \int_0^{(\alpha+r)} x^{a+r+s-2} \exp[-1/x] \, dx$$
$$= \frac{(-1)^a \, 2\pi i}{\Gamma(a)} F_3(-m,-n,m+a-1,n+a-1;a;1,1)$$

where F_3 is an Appell function of the third kind. See Exton [1] (page 24) for example.

The Appell function may be re-arranged in the form

(10)
$$F = \sum_{r=0}^{m} \frac{(-m)_r (m+\alpha-1)_r}{(a)_r r!} 2F_1(-n, n+\alpha-1; \alpha+r; 1)$$

and on applying Vandermonde's theorem, the inner Gauss polynomial becomes

$$\frac{(r-n+1)_n}{(r+a)_n}.$$

If r < n, the numerator of this expression vanishes and since r runs from 0 to m, it follows that F vanishes when m < n. By symmetry, it will be seen that F also vanishes when n < m. If m = n, then all the terms of the series (10) vanish except that for which r = n and we have

(12)
$$F = \frac{(-n)_n (n+a-1)_n}{(a)_n (n+a)_n} = \frac{(-1)^n n! (n+a-1)}{(a)_n (2n+a-1)}.$$

We may then write

(13)
$$I_{m,n} = \int_{0}^{(\alpha+1)} x^{\alpha-2} \exp\left[-\frac{1}{x}\right] y_m(x; a, 1) y_n(x; a, 1) dx$$
$$= \frac{(-1)^{\alpha+n} n! (n+\alpha-1) (2\pi i) \delta_{m,n}}{\Gamma(\alpha+n) (2n+\alpha-1)},$$

which is a new orthogonality relation involving the Bessel polynomials, with an elementary weight function.

Similarly, if, for convergence, Re(a) < 1 - m - n, we have

(14)
$$J_{m,n} = \int_{0}^{\infty} x^{a-2} \exp\left[-\frac{1}{x}\right] y_{m}(x; a, 1) y_{n}(x; a, 1) dx$$
$$= \frac{(-1)^{n} n! (n+a-1)\pi}{\Gamma(a+n) (2n+a-1) \sin(\pi a)} \delta_{m,n}.$$

References

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Abstract

A new orthogonality relation for the Bessel polynomials is given with an elementary weight function.

