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## On a Theorem of Louis de Branges. (\*\*)

1. - Louis de Branges [1] proved the following theorem for Hankel transform of order  $\nu$ .

Theorem. A necessary and sufficient condition for the functions g(x),  $f(x) \in L_2$  to be the Hankel transforms of one another is that the equation

(1.1) 
$$\int_{0}^{\infty} f(t) (xt)^{\nu+(1/2)} e^{-x^{2} t^{2}/2} dt = \int_{0}^{\infty} g(t) (x^{-1} t)^{\nu+(1/2)} e^{-t^{2}/(2x^{2})} x^{-1} dt ,$$

holds for all x > 0,  $\nu > -1$ , where  $x^{\nu + (1/2)} e^{-x^2/2}$  is defined as a fundamental self-reciprocal function for the Hankel transformation of order  $\nu$ .

In this paper we generalise the above theorem for any transformation, the kernel function of which is a symmetrical FOURIER kernel.

2. - Now we give certain results used in the following sections.

Let  $s = \sigma + it$  be a complex variable. Following Titchmarsh ([6], p. 252) the author [4] has established the following result.

A necessary and sufficient condition that a function f(x),  $\in A(\alpha, a)$  ([6], p. 252) should be its own k-transform, where k(x), the kernel function of the transform is such that its Mellin transform, K(s) is O(1),  $K_1(s)$  is  $O(e^{\lambda|t|})$  and  $K_1(s)$  satisfies the relation

(2.1) 
$$K(s) = K_1(s)/K_1(1-s)$$

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<sup>(\*\*)</sup> Ricevuto: 15-X-1968.

is that f(x) should be of the form

(2.2) 
$$f(x) = \frac{1}{2 \pi i} \int_{s-s}^{s+i\infty} K_1(s) \ \psi(s) \ x^{-s} \, \mathrm{d}s,$$

with

$$(2.3) \psi(s) = \psi(1-s),$$

where  $\psi(s)$  and  $K_1(s)$  are regular in the strip

(2.4) 
$$a < \sigma < 1 - a, \quad a < 1/2,$$

 $\psi(s)$  is  $O(e^{(-\lambda - \alpha + \eta)|t|})$  for every positive  $\eta$  and uniformly in the strip (2.4) and e in any value of  $\sigma$  in the strip (2.4).

In the equation

(2.5) 
$$g(x) = \int_{0}^{\infty} f(g) \ k(xy) \ dy$$

we can conclude from  $L_2$  theory following Titchmarsh ([6], p. 221) that g(x),  $f(x) \in L_2$  if k(x) is a symmetrical Fourier kernel.

Denoting the Mellin transform of f(x) by F(s), the Parseval relation ([3], p. 391) for Mellin trasform, with parameter x (x > 0) is given by

(2.6) 
$$\int_{0}^{\infty} f_{1}(ux) f_{2}(u) du = \frac{1}{2 \pi i} \int_{(1/2)-i\infty}^{(1/2)+i\infty} F_{1}(s) F_{2}(1-s) x^{-s} ds,$$

where  $f_1(x)$ ,  $f_2(x) \in L_2$  and  $F_1(s)$ ,  $F_2(s) \in L_2(\frac{1}{2} - i \infty, \frac{1}{2} + i \infty)$ .

We now extend the definition of fundamental self-reciprocal function for the Hankel transform of order  $\nu$  to k-transform.

If in (2.2)  $\psi(s) = 1$ , we have

(2.7) 
$$f(x) = k_f(x) = \frac{1}{2 \pi i} \int_{-\infty}^{c+i\infty} K_1(s) x^{-s} ds,$$

then  $k_{f}(x)$  is called the fundamental self-reciprocal function for the k-transform.

3. – Theorem. A necessary and sufficient condition for the functions g(x),  $f(x) \in A(\alpha, a)$  and  $L_2$ ,  $K(s) \in L_2(\frac{1}{2} - i\infty, \frac{1}{2} + i\infty)$ , and  $K_1(s)$  satisfies (2.1) and  $\in L_2(\frac{1}{2} - i\infty, \frac{1}{2} + i\infty)$  to be the k-transform of one another is that the equation

(3.1) 
$$\int_{0}^{\infty} f(t) k_{f}(xt) dt = \int_{0}^{\infty} g(t) k_{f}(t/x) \frac{dt}{x}$$

holds good for all x > 0, where  $k_f(x)$  is the fundamental self-reciprocal function for the k-transform.

Proof. Condition is necessary. We have

(3.2) 
$$g(x) = \int_{0}^{\infty} f(y) \ k(xy) \ dy.$$

As  $f(y) \in L_2$  and  $K(s) \in L_2(\frac{1}{2} - i\infty, \frac{1}{2} + i\infty)$  we can use Parseval relation (2.6) on the R.H.S. of (3.2) thus we have

(3.3) 
$$g(x) = \frac{1}{2\pi i} \int_{(1/2)-i\infty}^{(1/2)+i\infty} K(s) F(1-s) x^{-s} ds, \qquad x > 0.$$

As  $g(x) \in L_2$  the Mellin transform of g(x) exists and (3.3) reduces to

$$(3.4) G(s) = F(1-s) K(s),$$

using (2.1) in (3.4) we have

$$(3.5) G(s) K_1(1-s) = F(1-s) K_1(s).$$

As  $K_1(s) \in L_2(\frac{1}{2} - i\infty, \frac{1}{2} + i\infty)$  so taking the inverse Mellin transform of both sides of (3.5), we have

$$\frac{1}{2\pi i} \int\limits_{(1/2)-i\infty}^{(1/2)+i\infty} G(s) \ K_1(1-s) \ x^{-s} \ \mathrm{d}s = \frac{1}{2\pi i} \int\limits_{(1/2)-i\infty}^{(1/2)+i\infty} K_1(s) \ F(1-s) \ x^{-s} \ \mathrm{d}s \ .$$

Using the Parseval relation (2.6) and (2.7), we get

(3.6) 
$$\int_{0}^{\infty} g(xu) k_{f}(u) du = \int_{0}^{\infty} k_{f}(tx) f(t) dt.$$

Replacing xu by t in left hand side of (3.6), we have

(3.7) 
$$\int_{0}^{\infty} g(t) \ k_{f}(t/x) \ x^{-1} \ dt = \int_{0}^{\infty} f(t) \ k_{f}(xt) \ dt.$$

Condition is sufficient. Retracing the steps from (3.7) to (3.2) it can be shown that g(x) is the k-transform of f(x). To show that f(x) is the k-transform of g(x) we retrace steps from (3.7) to (3.4), replace s by 1-s to obtain

(3.8) 
$$G(1-s) K_1(s) = F(s) K_1(1-s).$$

Moving on similar lines from (3.4) to (3.2), we get

$$f(x) = \int_{0}^{\infty} g(y) \ k(xy) \ \mathrm{d}y$$

instead of (3.2). This concludes sufficiency.

4. - By selecting different kernel functions we get relations for different transforms. We mention few cases,

Corollary 1. If

$$k(x) = x^{1/2} J_{\nu}(x), \qquad \nu > -1,$$

then

$$egin{align} K(s) &= 2^{s-(1/2)} \, arGammaigg(rac{1}{4} + rac{v}{2} + rac{s}{2}igg) \, igg/ \, arGammaigg(rac{3}{4} + rac{v}{2} - rac{s}{2}igg) \, , \ & K_1(s) \, = 2^{s/2} \, arGammaigg(rac{1}{4} + rac{v}{2} + rac{s}{2}igg) \, . \end{split}$$

giving fundamental self-reciprocal function  $x^{\nu+(1/2)}e^{-x^2/2}$ .

Using the theorem, equation (3.1) reduces to (1.1) which is a known result.

Corollary 2. If  $k(x) = \omega_{\mu,\nu}(x)$   $(\mu > -1, \nu > -1)$  the kernel introduced by Watson [7], then

$$K_1(s) = 2^s \ arGamma\left(rac{1}{4} + rac{v}{2} + rac{s}{2}
ight) \ arGamma\left(rac{1}{4} + rac{\mu}{2} + rac{s}{2}
ight),$$

and the fundamental self-reciprocal function for Watson transform is

$$2 G_{0,2}^{2,0} \left( \frac{x^2}{4} \middle| \frac{1}{2} + \frac{\mu}{2}, \frac{1}{4} + \frac{\nu}{2} \right).$$

Using the result ([2], p. 434) this becomes

$$2^{(3-\mu-\nu)/2} x^{(\mu+\nu+1)/2} K_{(\mu-\nu)/2}(x)$$

where  $K_{\alpha}(x)$  denotes the modified Bessel function of order  $\alpha$ . Using the theorem, the equation (3.1) assumes the form

$$(4.1) \quad \int\limits_0^\infty g(t) \ (tx^{-1})^{(\mu+\nu+1)/2} \ K_{\,(\mu-\nu)/2}(tx^{-1}) \ x^{-1} \ \mathrm{d}t = \int\limits_0^\infty f(t) \ (xt)^{(\mu+\nu+1)/2} \ K_{\,(\mu-\nu)/2}(xt) \ \mathrm{d}t \ .$$

Corollary 3. If  $k(x) = \chi_{\nu,k,m}(x)$  the kernel function studied by R: N ar a in ([5], p. 271), then

(4.2) 
$$K_1(s) = \frac{2^{s/2} \Gamma\left(\frac{1}{4} + \frac{\nu}{2} + \frac{s}{2}\right) \Gamma\left(\frac{1}{4} + \frac{\nu}{2} + 2m + \frac{s}{2}\right)}{\Gamma\left(\frac{3}{4} + \frac{\nu}{2} + m - k + \frac{s}{2}\right)},$$

$$\operatorname{Re} s \geqslant s_0 > 0 \; , \qquad \qquad \operatorname{Re} \; (\nu \; + 1 \; + 2m \; \pm \; 2m) > 0 \; ,$$

2m not an integer and the fundamental self-reciprocal function for  $\chi_{\nu,k,m}$ -transform is given by

$$G_{1,2}^{2}\left(\frac{x^{2}}{2}\left|\frac{\frac{3}{4}+\frac{\nu}{2}+m-k}{\frac{1}{4}+\frac{\nu}{2},\frac{1}{4}+\frac{\nu}{2}+2m}\right)\right), \qquad \text{Re}\left(\nu+1+2m\pm2m\right)>0.$$

Applying the theorem, the equation (3.1) becomes

(4.3) 
$$\int_{0}^{\infty} f(t) G_{1,2}^{2,0} \left( \frac{t^{2} x^{2}}{2} \left| \frac{\frac{3}{4} + \frac{\nu}{2} + m - k}{\frac{1}{4} + \frac{\nu}{2}, \frac{1}{4} + \frac{\nu}{2} + 2m} \right) dt =$$

$$= \int_{0}^{\infty} g(t) G_{1,2}^{2,0} \left( \frac{t^{2}}{2 x^{2}} \left| \frac{\frac{3}{4} + \frac{\nu}{2} + m - k}{\frac{1}{4} + \frac{\nu}{2}, \frac{1}{4} + \frac{\nu}{2} + 2m} \right) x^{-1} dt.$$

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## Abstract.

In this paper we have generalised the theorem of Hankel transform given by Louis de Branges.

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